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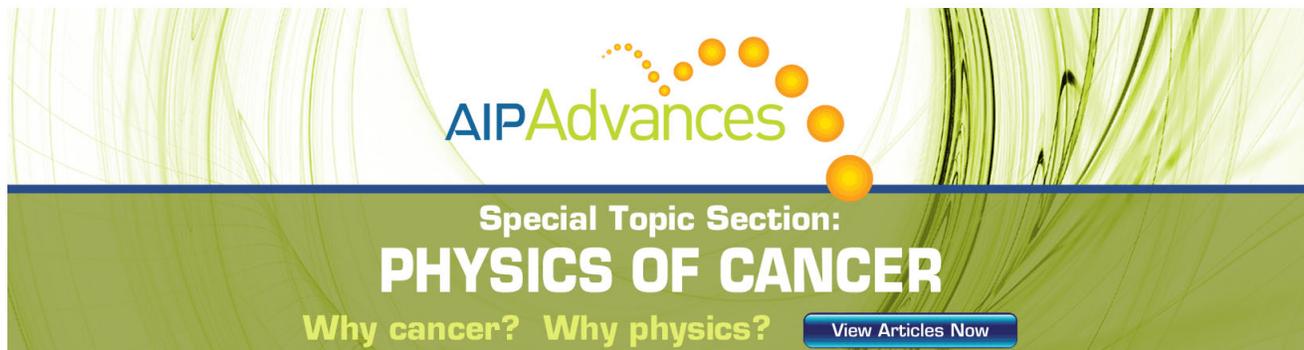
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Assessment of density functional theory for iron(II) molecules across the spin-crossover transition

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Octahedral Fe²⁺ molecules are particularly interesting as they often exhibit a spin-crossover transition. In spite of the many efforts aimed at assessing the performances of density functional theory for such systems, an exchange-correlation functional able to account accurately for the energetic of the various possible spin-states has not been identified yet. Here, we critically discuss the issues related to the theoretical description of this class of molecules from first principles. In particular, we present a comparison between different density functionals for four ions, namely, [Fe(H₂O)₆]²⁺, [Fe(NH₃)₆]²⁺, [Fe(NCH)₆]²⁺, and [Fe(CO)₆]²⁺. These are characterized by different ligand-field splittings and ground state spin multiplicities. Since no experimental data are available for the gas phase, the density functional theory results are benchmarked against those obtained with diffusion Monte Carlo, one of the most accurate methods available to compute ground state total energies of quantum systems. On the one hand, we show that most of the functionals considered provide a good description of the geometry and of the shape of the potential energy surfaces. On the other hand, the same functionals fail badly in predicting the energy differences between the various spin states. In the case of [Fe(H₂O)₆]²⁺, [Fe(NH₃)₆]²⁺, [Fe(NCH)₆]²⁺, this failure is related to the drastic underestimation of the exchange energy. Therefore, quite accurate results can be achieved with hybrid functionals including about 50% of Hartree-Fock exchange. In contrast, in the case of [Fe(CO)₆]²⁺, the failure is likely to be caused by the multiconfigurational character of the ground state wave-function and no suitable exchange and correlation functional has been identified. © 2012 American Institute of Physics. [<http://dx.doi.org/10.1063/1.4752411>]

I. INTRODUCTION

Among the transition metal complexes, octahedral 3d⁶ Fe²⁺ molecules are systems of particular interest. In fact they often undergo the spin-crossover (SC) transition.¹ In their most common form, the molecules' ground state is a spin singlet (¹A_{1g} symmetry in group theory notation), with the Fe six 3d electrons paired up in the t_{2g}π* antibonding orbitals. Upon increasing the temperature, the high spin (HS) quintet state (⁵T_{2g}), in which two electrons are promoted from the t_{2g}π* to the e_gσ* orbitals, becomes thermodynamically more stable (see Fig. 1 for the molecular orbital diagram). The SC transition is entropy driven and it is regulated by the difference in the Gibbs free energy of the HS ⁵T_{2g} and low spin (LS) ¹A_{1g} states,

$$\Delta G = G_{\text{HS}} - G_{\text{LS}} = \Delta H - T\Delta S. \quad (1)$$

Here, $\Delta H = H_{\text{HS}} - H_{\text{LS}}$ and $\Delta S = S_{\text{HS}} - S_{\text{LS}}$ indicate, respectively, the enthalpy and the entropy variation (note that for $\Delta G > 0$ the LS configuration is more thermodynamically stable than the HS one). For SC molecules $\Delta H > 0$, but also $S_{\text{HS}} > S_{\text{LS}}$, i.e., $\Delta S > 0$. Hence, for large enough temperatures ($T > T_c = \Delta H/\Delta S$), the entropic term dominates over the enthalpic one and the molecules transit from the LS to the HS configuration. There are two contributions to the entropy: the first is provided by the spin and the second by

the molecule vibrations. In fact, the transfer of two electrons to the e_gσ* orbitals, which are more antibonding than the t_{2g}π*, weakens the chemical bond and produces a breathing of the metal ion coordination sphere. This results in a softening of the phonon modes and then an increase of the vibronic entropy.

The SC transition is usually reported either for molecules in solution or in single crystals and, depending on the strength and on the origin of the inter-molecular interactions, it can have cooperative nature and present a thermal hysteresis loop. Interestingly, the transition can be also induced by illumination. This phenomenon is called light-induced-excited-spin-state-trapping effect (LIESST) and it is explained through a mechanism involving several excited states.² The lifetime of the metastable HS state is usually very long at low temperature as the relaxation to the ground state is due to the weak electronic coupling between these states.³ The opposite process, in which molecules populating the metastable HS state are brought back to their ground state, is also possible and it is called reverse LIESST effect.

Although, SC molecules have been traditionally studied for possible applications in storage and sensor devices,^{1,4,5} they have recently emerged as promising materials for molecular spintronics.⁶⁻⁸ In fact, the electronic transport through these molecules has been predicted^{9,10} and

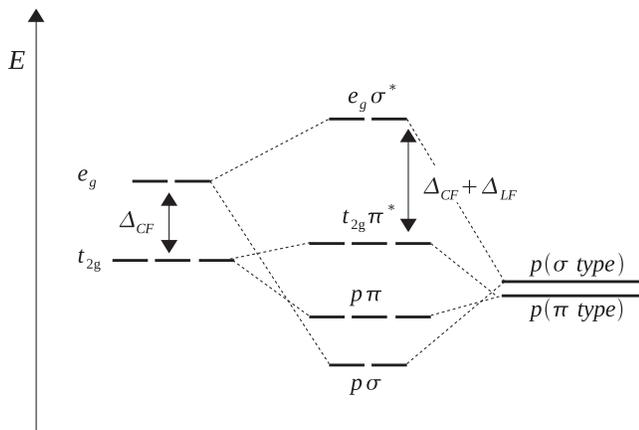


FIG. 1. Energy level diagram for an octahedrally coordinated transition metal (TM) ion. In crystal field theory, the $3d$ orbitals of the TM ion have an energy split Δ_{CF} due to the electrostatic interaction with the ligands. In contrast, in ligand field theory the $3d$ orbitals of the TM ion form covalent bonds with the ligands. In this diagram, we assume that each ligand contributes three p -orbitals, one with the positive lobe pointing toward the TM ion, σ -type, and two with the lobes perpendicular to it, π -type, (note that the σ - and π -type p orbitals are degenerate but here they are plotted slightly separated for better display). The π -type p orbitals couple with the TM t_{2g} states, while the σ -type p orbitals couple with the e_g . Since the π interaction is weaker than the σ one, the antibonding $t_{2g}\pi^*$ orbitals lie lower in energy than the $e_g\sigma^*$ ones. The energy splitting between the $t_{2g}\pi^*$ and the $e_g\sigma^*$ orbitals is indicated by $\Delta_{LF} + \Delta_{CF}$.

experimentally reported^{11–14} to depend strongly on their spin state.

Given such renewed interest in spin crossover compounds there is also a growing fundamental effort in modeling their properties. In this respect, one aims at using an electronic structure theory, which is at the same time accurate and scalable. Accuracy is needed for reliable predictions of the spin crossover temperature, while scalability is required by the size of the typical molecules. This becomes particularly crucial for molecules in crystals and when deposited on metallic surfaces, since the typical simulation cells are large. Density functional theory (DFT) is in principle both scalable and accurate, but to date it is completely unclear how it does perform relatively to this problem.

In this paper, we investigate the performances of several commonly used exchange–correlation density functionals for predicting the relative energy of the HS $^5T_{2g}$ and the LS $^1A_{1g}$ state of four model Fe^{2+} ions, namely, $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$, $[\text{Fe}(\text{NH}_3)_6]^{2+}$, $[\text{Fe}(\text{NCH})_6]^{2+}$, and $[\text{Fe}(\text{CO})_6]^{2+}$. We then compare our DFT results to those of diffusion Monte Carlo (DMC).^{15–17} Since the experimental data for Fe^{2+} complexes are difficult to compare with a microscopic theory, our work provides a quite informative benchmark for the theory itself.

The paper is organized as follows. In Sec. II, we present an overview of the problem and of the various known shortcomings of DFT, and we discuss critically which elements one has to consider when comparing electronic structure data to experiments. Then we will provide some computational details and move to the results. First, we will discuss our DFT calculations for the four different ions and then we will compare them with the DMC ones. Finally, we will conclude.

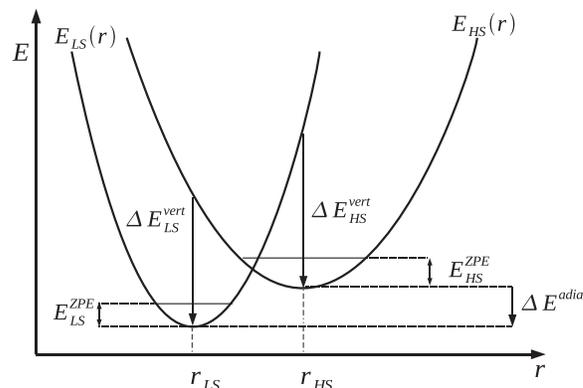


FIG. 2. Potential energy surface of the HS and LS state of a SC molecule. The collective coordinate r represents all the $3N$ nuclear coordinates of the molecule. The zero-point phonon energies for the HS and LS state, $E_{\text{HS}}^{\text{ZPE}}$ and $E_{\text{LS}}^{\text{ZPE}}$, the adiabatic energy gap, ΔE^{adia} , and the vertical energy gaps, $\Delta E_{\text{LS}}^{\text{vert}} = \Delta E^{\text{vert}}(r_{\text{LS}})$ and $\Delta E_{\text{HS}}^{\text{vert}} = \Delta E^{\text{vert}}(r_{\text{HS}})$ are indicated.

II. DENSITY FUNCTIONAL THEORY FOR SPIN-CROSSOVER MOLECULES: STATE OF ART

When one considers a single molecule in vacuum at zero temperature, ΔG coincides with the internal energy difference, which in the adiabatic approximation is

$$\Delta E = \Delta E^{\text{adia}} + \Delta E^{\text{ZPE}}. \quad (2)$$

Here, $\Delta E^{\text{ZPE}} = E_{\text{HS}}^{\text{ZPE}} - E_{\text{LS}}^{\text{ZPE}}$ and $E_{\text{HS}}^{\text{ZPE}}$ ($E_{\text{LS}}^{\text{ZPE}}$) is the zero-point phonon energy of the HS (LS) state, while

$$\Delta E^{\text{adia}} = E_{\text{HS}}(r_{\text{HS}}) - E_{\text{LS}}(r_{\text{LS}}) \quad (3)$$

is the adiabatic energy difference (often called “adiabatic energy gap”). The collective coordinate r represents the $3N$ nuclear coordinates of the molecule and the energies $E_{\text{HS}}(r)$ and $E_{\text{LS}}(r)$ define the potential energy surfaces (PESs), respectively, of the HS and LS state (see Fig. 2). In addition to the adiabatic energy gap we can also define the vertical energy difference (“vertical energy gap”)¹⁸

$$\Delta E^{\text{vert}}(r) = E_{\text{HS}}(r) - E_{\text{LS}}(r) \quad (4)$$

and the difference of vertical energy gaps (DOG)¹⁹

$$\text{DOG}(r_{\text{HS}}, r_{\text{LS}}) = \Delta E^{\text{vert}}(r_{\text{HS}}) - \Delta E^{\text{vert}}(r_{\text{LS}}). \quad (5)$$

All of these quantities can be computed by using *ab initio* electronic structure methods. As we mentioned in the Introduction, DFT is always the preferred one. In fact SC molecules are composed of, at least, 50 atoms and a good balance between expected accuracy and computational cost is required. However, there are many issues connected to the DFT description of SC molecules, which either have not been properly addressed or have not found any satisfactory solution yet. Here, we list some of them.

A. The functional dilemma

For each Fe^{2+} molecules, in general, every exchange–correlation functional returns a very different adiabatic energy gap (see, for example, Refs. 19 and 20). These differences can be as large as few eV and different functionals do

not sometimes even predict the same ΔE^{adia} sign. In a nutshell, no agreement around which functional performs best has been reached so far (the discussion below will explain how problematic is a direct assessment of the DFT results through a comparison with the experimental data). However, some general trends, which relate functionals belonging to the same “class” (or the same “rung” if we refer to the “Jacob’s ladder”²¹ classification scheme), have been identified.

1. The local density approximation (LDA) (first rung) tends to stabilize the LS state

This is due to the underestimation of the exchange energy.²² In particular, the exchange keeps electrons of like-spin apart so that their Coulomb repulsion is reduced. Therefore, the exchange underestimation is accompanied by the overestimation of the Coulomb energy for two electrons of equal spin. This, in turns, leads to the stabilization of the LS state at the expense of the HS state.

2. The generalized gradient approximation (GGA) and the meta-GGAs (second and third rungs) give results that depend on the specific compound and on the exact DFT conditions that each functional satisfies

“Traditional” GGAs,²³ such as the Perdew-Burke-Ernzerhof (PBE)²⁴ and the BLYP, which combines the Becke exchange²⁵ with the Lee-Yang-Parr (LYP) correlation,²⁶ reduce only slightly the LDA over-stabilization of the LS state. Therefore, they do not represent a drastic improvement. In contrast, some more recent GGA functionals, such as the OLYP (combining the Handy and Cohen’s OPTX exchange²⁷ with the LYP correlation), have been claimed to perform rather well.^{28,29,31} Among the meta-GGAs, the Van Voorhis-Scuseria exchange-correlation functional³⁰ was tested by Ganzenmüller *et al.*,³¹ who concluded that it provided quite accurate results for single-point calculations, while it predicted artificially twisted structures.

3. Hybrid functionals (forth rung) tend to favor the HS with respect to LS state

Reiher and co-workers^{32,33} argued that the amount of Hartree-Fock (HF) exchange in many hybrid functionals is too large to predict correct energy gaps. They then proposed a re-parametrization of the B3LYP functional,³⁴ called B3LYP*, which includes only 15% of HF exchange (in contrast to the standard B3LYP 20%). Although B3LYP* is believed to give satisfactory results for several SC complexes, some studies suggested that a further reduction of the HF exchange could be needed in order to describe other Fe²⁺ compounds.^{31,32} In contrast, the amount of HF exchange in B3LYP was judged insufficient for the small ions [Fe(H₂O)₆]²⁺ and [Fe(NH₃)₆]²⁺. For these systems, it has been claimed that PBE0,^{35,36} which includes up to 25% of HF exchange, gives more satisfactory results.^{28,29} In practice, for each compound, the amount of HF exchange can be varied to fit the desired values for the gaps but no “universally good choice” has been identified so far. The dependence of

the adiabatic energy gap on the amount of HF and either of LDA (Slater) or GGA (B88 and OPTX) exchange is explained very clearly in the work of Ganzenmüller *et al.*³¹ Finally, it is important to remark that, even when a certain hybrid functional is found to return satisfactory energy gaps, it might not be the optimal functional to describe other properties of the molecule.

B. Basis set

Density functional theory calculations for SC molecules are usually performed by using Quantum Chemistry codes, where the wave-function is expanded over either Gaussian (GTOs) or Slater-type orbitals (STOs).³⁷ In many cases, the values of the energy gaps depend substantially on the choice of the basis sets and on the spatial extension of the local orbitals.³⁸ Although this is a less severe problem compared to that of identifying the exchange-correlation functional, it must be kept into consideration. In principle, the use of plane-waves (PW) basis sets,³⁹ instead of GTOs and STOs, could be a solution, but in practice plane-wave calculations are computationally expensive because of the need of satisfying periodic boundary conditions.

A large number of plane-waves is usually needed as the electronic density is concentrated in a small fraction of the supercell volume. Furthermore, very large supercells are typically required. This is due to the fact that SC complexes are often 2+ ions. Therefore, a negatively charged background must be introduced in order to preserve the overall charge neutrality so that the total energy remains bound. The energy calculated in this way approaches then the one for an isolated system only in the limit of large supercell and, unfortunately, such convergence is slow. Although corrections to the expression of the computed energy have been proposed,^{40,41} this effect can be properly accounted for only by considering large supercells and by performing multiple calculations for supercells of different sizes.

C. Geometry optimization

Each class (rung) of exchange-correlation functionals returns different metal-ligands bond-lengths. Usually, LDA gives shorter bonds than hybrids functionals, while standard GGA results are in between these two extremes.^{42,43} Although these differences are usually quite small, less than 0.1 Å against an average bond-length of about 2 Å, they might strongly affect the electronic properties. Indeed, a very delicate balance between ligand-field splitting and Hund’s coupling establishes the spin state (see Fig. 1). Then, small errors in the predicted geometries can drastically alter this balance.

Unfortunately, the quality of the DFT relaxation cannot be easily assessed through a comparison with available experiments. In fact, while DFT calculations are usually carried out for molecules in the gas phase, the experimental geometries are obtained through x-ray measurements for crystals.⁴⁴ As far as we know, no detailed DFT studies about SC molecules in the condensed phase have been published so far. Furthermore, such a study must face the additional difficulty of the

need of accounting for inter-molecular interactions. These have usually dispersive nature and, therefore, they are either not described or badly described by most of the popular functionals.⁴⁵

D. Electrostatic contributions to the total energy

Spin-crossover molecules, which are usually in the 2+ charging state, are surrounded by counter-ions (for example, $[\text{PF}_6]^{2-}$ or $[\text{BF}_4]^{2-}$). Because of the electrostatic potential generated by such counter-ions and by the other molecules in the crystal, the total energy of a SC complex in the condensed phase differs from that in the gas phase. Furthermore, since at the SC phase transition there is a charge redistribution over the individual molecules and a lattice expansion, such electrostatic potential does not induce a simple rigid shift of the minima of the HS and LS PESs (compared to those calculated for the gas phase). The energy gap then turns out to be different for the same molecule in different phases.^{46,47} Unfortunately, this effect is always neglected by DFT calculations, which aim at assessing the performances of DFT by using experimental data.

E. Finite-temperature effects

Special care should be taken in order to include finite-temperature effects when comparing DFT to experiments. In fact, at finite temperature instead of the adiabatic energy gap, the Gibbs free energy, Eq. (1), must be computed. So far, this has been attempted only by Ganzenmüller *et al.*³¹ However, unfortunately, their calculations did not fully account for either the electrostatic contribution to the total energy or the effect of the periodic lattice on the molecular structure.

This list clearly emphasizes how the main handicap in the theoretical description of SC complexes is related to the difficulties in assessing the performances of any given exchange-correlation functional. Any benchmark involving a comparison against experimental results is fated to fail, unless vibrational, environmental, crystallographic, and finite-temperature effects are properly accounted for. However, this task is at present too demanding to be practically achievable. In contrast, as pointed out by Fouqueau *et al.*,^{28,29} the current best strategy consists in providing benchmark values for various interesting quantities through highly accurate *ab initio* methods. These can then be compared with the DFT results in order to identify which functional performs better.

In some interesting works,^{14,28,29,48,49} wave-function based methods were considered (see discussion below). However, unfortunately, the authors themselves admitted that their results were plagued by a number of systematic errors. These were ascribed to the too small basis set used for the Fe^{2+} ion and to the fact that such computational methods describe exactly only for static electronic correlation, but do not include the dynamic one, which, however, can be partially included by perturbation theory.

Here, we have chosen to employ DMC as benchmark electronic structure theory. Instead of describing the correlation energy by an approximate exchange-correlation

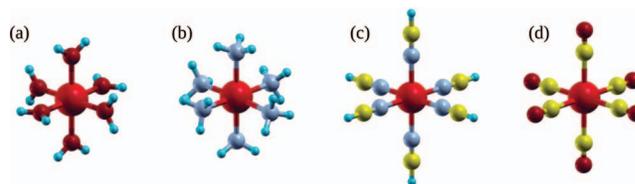


FIG. 3. The cations investigated in this work $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$ (a), $[\text{Fe}(\text{NH}_3)_6]^{2+}$ (b), $[\text{Fe}(\text{NCH})_6]^{2+}$ (c), and $[\text{Fe}(\text{CO})_6]^{2+}$ (d). Color code: C = yellow, O = red (small sphere), Fe = red (large sphere), N = grey, H = blue.

energy functional, DMC is based on a stochastic sampling of the many-electron wave-function, thus treating electronic correlation in a truly many-electron fashion. Although, at the moment, there are only a few attempts to apply DMC to transition metal complexes (mainly bi-atomic transition metal oxide molecules),⁵⁰ so that a thorough assessment of the method is not available, we believe that the DMC predictions still represent a drastic improvement over the DFT results. Furthermore, DMC may be complementary to the wave-function based methods as the dynamic correlation is usually well described by using an appropriate Jastrow factor (see Sec. III). Unfortunately, however, the final accuracy of the calculation may be affected by the fixed-node approximation. Here, we will briefly mention some issues related to such approximation, to which we will dedicate a more detailed study in the future.

In this work, we compare systematically several DMC results to those obtained with DFT for a few selected Fe^{2+} complexes. Unfortunately, such a systematic investigation requires a large use of computational resources and, therefore, it cannot be carried out for molecules composed of tens of atoms (such as the most typical SC complexes). We have then focused our attention on the following ions: $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$, $[\text{Fe}(\text{NH}_3)_6]^{2+}$, $[\text{Fe}(\text{NCH})_6]^{2+}$, and $[\text{Fe}(\text{CO})_6]^{2+}$, which are shown in Fig. 3. These are small enough to allow several DMC calculations to be performed at a reasonable computational cost. Furthermore, and more importantly, the study of their electronic structure presents all the problems mentioned above so that our main conclusions can be extended to large SC molecules as well. Finally, according to the spectrochemical series,⁵¹ these ions have a different ligand field splitting. The one of $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$ is the smallest and that of $[\text{Fe}(\text{CO})_6]^{2+}$ the largest. We then expect that $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$, $[\text{Fe}(\text{NH}_3)_6]^{2+}$, and $[\text{Fe}(\text{NCH})_6]^{2+}$ have a HS ground state, while $[\text{Fe}(\text{CO})_6]^{2+}$ a LS one. Therefore, our study scans through systems of different ground state spin multiplicity, it reveals general trends and it points to the systematic errors of each DFT functional.

III. COMPUTATIONAL DETAILS

The DFT calculations are performed with the NWCHEM package.⁵² We use several functionals belonging to different “classes”: (1) the default LDA of Vosko-Wilk-Nussair,⁵³ (2) the GGA BP86 functional, which combines the Becke88 exchange functional²⁵ with the Perdew86 correlation one,⁵⁴ and (3) the hybrid functionals B3LYP,³⁴ PBE0,^{35,36} and the

Becke “half and half” (Becke-HH).⁵⁵ These include, respectively, 20%, 25%, and 50% of HF exchange (note that, in the NWCHEM package, only an approximate version of the HH-Becke functional is currently implemented⁵⁶).

We have chosen three different basis sets: (1) 6-31G* (basis set called A), (2) the LANL2DZ (8s,5p,5d)/[3s,3p,2d] basis set and pseudopotential⁵⁷ for Fe combined with the basis set 6-31++G** for all other atoms (basis set B), and (3) the triple-zeta polarized basis set of Ahlrichs⁵⁸ (basis set C) which is a (11s,5p,1d)/[5s,3p,1d] basis for C, N, and O and (17s,11p,6d)/[6s,4p,3d] basis for Fe. Although in some cases, we have checked the convergence of the results with respect to the quality of the basis set (for example, by using the Ahlrichs basis set the results are converged already at the double zeta level), we present here only the results for the basis sets A, B, and C as these were used in similar studies (see, for example Refs. 19, 28, and 29). Our DFT results can then be directly compared to those of other authors.

Geometry optimizations are performed both with and without specifying the molecule point group. The geometries optimized with these two strategies usually return consistent results with bond-lengths differences, which are within ± 0.005 Å. We always check that the phonon frequencies are all real so that the final geometries correspond to stable energy minima.

The adiabatic energy gaps are computed at the geometry obtained with the same basis set. As we will mention in Secs. IV and V we could not establish which basis set returns the lowest energy structure. However, for a given functional, we find that the adiabatic energy gaps change at most by about 0.1 eV when a different basis set (and the corresponding optimized geometry) is considered. This change is much smaller than the difference between the gaps predicted by using the different functionals. As our aim here is that of comparing the performances of different density functionals with those of DMC, such minor uncertainty over the choice of basis set and the associated geometry plays only a secondary role on our conclusions.

The DMC calculations are performed by using the CASINO code.¹⁶ The trial wave-functions are of Slater-Jastrow type

$$\Psi_T(\mathbf{R}) = e^J D^\uparrow D^\downarrow, \quad (6)$$

where D^\uparrow and D^\downarrow are the Slater determinants of the up and down single-particle orbitals. The quantity e^J is the Jastrow factor with the exponential J including a sum of one-body (electron-nucleus), two-body (electron-electron), and three-body (electron-electron-nucleus) terms. These are functions, which depend on the electron-nucleus, electron-electron, and electron-electron-nucleus separations, and which satisfy the cusp conditions.⁵⁹ The parameters in the Jastrow factor are optimized by minimizing the variational Monte Carlo energy variance. We note that, by requiring Ψ_T to have A_{1g} (T_{2g}) symmetry and spin equal to 0 (2), DMC simulations effectively return an upper bound to the total electronic energy of the ${}^1A_{1g}$ state (the ${}^5T_{2g}$ state).

Slater-Jastrow trial wave-functions describe very well dynamic correlation effects. In contrast, in order to account accurately for static correlation, a multi-determinant trial wave-

function may be more suitable. However, as we will discuss below, the ${}^5T_{2g}$ and ${}^1A_{1g}$ states have a single-configuration character so that a Slater-Jastrow trial wave-function is appropriate. Besides, simulations employing a multi-determinant trial wave-function are in general not feasible for large systems. In fact a large number of determinants is required for an accurate estimate of the static correlation energy (the number of determinants grows at least polynomially, if not even exponentially with the system size).

The imaginary time evolution of the Schrödinger equation has been performed with the usual short time approximation and time-steps of various sizes are considered (typically $\Delta\tau = 0.0125, 0.005, 0.001$ a.u.). Except for few cases, the energy differences are usually found to be converged with respect to the time-step errors already for $\Delta\tau = 0.0125$ a.u. Calculations are performed by using Dirac-Fock pseudopotentials^{60,61} with the “potential localization approximation” (PLA).⁶² For $[\text{Fe}(\text{NCH})_6]^{2+}$, DMC simulations with this approximation are found to be unstable as the number of walkers “explodes.” Therefore, we have used instead the “T-move” scheme,^{63,64} which eliminates the need of the PLA and treats the non-local part of the pseudopotential in a way consistent with the variational principle. The simulations then become more stable. The single-particle orbitals of the trial wave-function are obtained through (LDA) DFT calculations performed with the PW code QUANTUM ESPRESSO.⁶⁵ The same pseudopotentials used for the DMC calculations are employed. The PW cutoff is fixed at 300 Ry and the PW are re-expanded in terms of B-splines.⁶⁶ The B-spline grid spacing is $a = \pi/G_{\max}$, where G_{\max} is the length of the largest vector employed in the PW calculations. Periodic boundary conditions are employed for the PW-DFT calculations and supercells as large as 40 Å are considered. In contrast, no periodic boundary conditions are imposed for the DMC simulations.

The DMC calculations are performed for the molecular geometries previously optimized by DFT. Therefore, we can compare the DMC energies of molecular structures obtained by employing different functionals and basis sets. However, for each system, these energy differences are often smaller than the computed error bars. Only LDA systematically returns molecular geometries with much higher DMC energies than those obtained by using either GGA or hybrid functionals.

IV. DFT RESULTS

A. $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$

The lower energy geometry of $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$ is found to have C_i symmetry for both BP86 and hybrid functionals. This is consistent with the results of Pierloot and Vancoillie.⁶⁷ In contrast, with LDA, we were able to obtain relaxed atomic positions for both the HS and LS states only by using the basis set A and without specifying the molecule point group.

As expected from our introductory discussion, the molecule in the LS state has metal-ligand bond-lengths shorter than those of the molecule in HS state (by about 7%). However, the details of the geometry depend on both the

TABLE I. Bond-lengths of $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$ in the HS and LS state, as calculated with various functionals and basis sets. Note that the LDA calculations for the HS state did not converge in the case of the basis sets B and C.

Functional	Basis set	d_{LS} (Å)	d_{HS} (Å)
LDA	A	1.917	2.052, 2.083, 2.057
BP86	A	1.985	2.152, 2.151, 2.111
BP86	B	2.02	2.174, 2.164, 2.132
BP86	C	2.01	2.161, 2.155, 2.125
B3LYP	A	2.005	2.152, 2.152, 2.111
B3LYP	B	2.003	2.146, 2.157, 2.112
B3LYP	C	2.029	2.172, 2.16, 2.137
PBE0	A	1.99	2.1, 2.145, 2.134
PBE0	B	2.013	2.168, 2.156, 2.129
PBE0	C	2.008	2.152, 2.147, 2.124
HH	B	2.010	2.168, 2.133, 2.132
HH	C	2.008	2.149, 2.131, 2.128

functional and the basis set. This can be clearly seen by inspecting Table I, which reports a full list of the calculated Fe–O bond-lengths for both the HS and LS states. On the one hand, LDA overbinds the molecule as compared to GGA and hybrids. On the other hand, the basis set A tends to shrink the Fe–O bond-lengths when compared to the basis set B and C. Although the choice of basis set does not affect the bond-lengths as drastically as the functional does, it still influences the geometry greatly. The calculations performed with the basis set A return a quite large inclination (about 5°) of the O–Fe–O axis with respect to the 90° angle it forms with the equatorial plane of the molecule. Furthermore, for basis set A and B, either the axial waters, which form the ligands, move “in” and the equatorial waters move “out” of their plane or vice versa. These results do not depend qualitatively on the functional. In contrast, the inclination of the O–Fe–O axis and the distortion of axial waters disappears when the calculations are carried out by using the basis set C.

Table II shows our calculated values for the adiabatic energy gaps, where a positive (negative) energy means that the LS state has lower (higher) energy than that of the HS one. For each functional, our results are always in very close agreement with those obtained by other authors^{29,48} (the results are

TABLE II. Adiabatic energy gap, ΔE^{adia} , for the cation $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$. The functional and the basis set used for the each calculation are indicated.

Functional	Basis set	ΔE^{adia} (cm^{-1})	ΔE^{adia} (eV)
LDA	A	–3986	–0.4942
BP86	A	–8989	–1.1145
BP86	B	–8381	–1.0391
BP86	C	–8400	–1.0415
B3LYP	A	–11 589	–1.4369
B3LYP	B	–11 027	–1.3672
B3LYP	C	–11 045	–1.3694
PBE0	A	–14 670	–1.8189
PBE0	B	–15 512	–1.9233
PBE0	C	–14 045	–1.7414
HH	C	–19 620	–2.4326
HH	C	–18 223	–2.2594

TABLE III. Bond-lengths of $[\text{Fe}(\text{NH}_3)_6]^{2+}$ in the HS and LS state, as calculated with various functionals and basis sets. Note that the LDA calculations for the HS state did not converge in the case of basis set B.

Functional	Basis set	d_{LS} (Å)	d_{HS} (Å)
LDA	A	1.942	2.188, 2.162, 2.160
LDA	C	1.995	2.204, 2.201, 2.214
BP86	A	2.026	2.267, 2.254, 2.253
BP86	B	2.078	2.30, 2.295, 2.274
BP86	C	2.085	2.279, 2.302, 2.289
B3LYP	A	2.076	2.281, 2.281, 2.275
B3LYP	B	2.114	2.315, 2.296, 2.294
B3LYP	C	2.122	2.32, 2.308, 2.283
PBE0	A	2.05	2.254, 2.256, 2.256
PBE0	B	2.082	2.292, 2.277, 2.272
PBE0	C	2.093	2.284, 2.294, 2.263
HH	C	2.11	2.296, 2.286, 2.266

presented in cm^{-1} as well as in eV in order to allow for a better comparison with the various values found in literature). Here, we can distinguish a clear trend, summarized by the series

$$-\Delta E^{\text{adia}}(\text{LDA}) < -\Delta E^{\text{adia}}(\text{GGA}) < -\Delta E^{\text{adia}}(\text{B3LYP}) \\ < -\Delta E^{\text{adia}}(\text{PBE0}) < -\Delta E^{\text{adia}}(\text{HH}). \quad (7)$$

This suggests that the calculated ΔE^{adia} is strictly related to the amount of HF exchange incorporated into the given functional. By increasing such contribution, we systematically stabilize the HS configuration with respect to the LS one.

B. $[\text{Fe}(\text{NH}_3)_6]^{2+}$

The optimized structure of the $[\text{Fe}(\text{NH}_3)_6]^{2+}$ ion, calculated either with BP86 or with hybrid functionals, has a D_3 symmetry for the LS state. This is further lowered to C_2 for the HS one. Our results are again consistent with those of Pierloot and Vancoillie.⁴⁸ Like in the case of $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$, we were not able to find the relaxed atomic geometry with LDA. Even when the geometry optimization procedure converges, like in the case of the basis set A and C, the minimum is found to be unstable. This is indicated by the negative eigenvalues of some phonon modes. Nevertheless, we report these results for completeness.

The molecule in LS state has shorter average Fe–ligand bond-lengths than the molecules in HS state (see Table III). In contrast to the case of $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$, $[\text{Fe}(\text{NH}_3)_6]^{2+}$ does not show any strong deviation of the N–Fe–N axis with respect to the axis normal to the equatorial plane for any combination of functionals and basis sets.

Table IV shows several values for the adiabatic energy gaps. Once again these can be ordered according to the series (7). Here, the LDA adiabatic energy gap indicates that $[\text{Fe}(\text{NH}_3)_6]^{2+}$ is LS. This result is even qualitatively incorrect as this cation is known to be stable in the HS state. BP86 also predicts the LS state to be the lowest in energy, although the value of ΔE^{adia} is very small and probably very sensible to the exact details of the calculation. In fact, in contrast to our results, which agree with those in Ref. 29, Pierloot and

TABLE IV. Adiabatic energy gap, ΔE^{adia} , for the cation $[\text{Fe}(\text{NH}_3)_6]^{2+}$. The functional and the basis set used for each calculation are indicated.

Functional	Basis set	ΔE^{adia} (cm^{-1})	ΔE^{adia} (eV)
LDA	A	8937	1.1081
LDA	C	7746	0.9605
BP86	A	195	0.0242
BP86	B	708	0.0878
BP86	C	672	0.0834
B3LYP	A	-5312	-0.6586
B3LYP	B	-4007	-0.4969
B3LYP	C	-4738	-0.5874
PBE0	A	-7695	-0.9541
PBE0	B	-7665	-0.9504
PBE0	C	-7117	-0.8825
HH	C	-13556	-1.68077

Vancoillie⁴⁸ obtained a negative value equal to about -0.2 eV. Finally, hybrid functionals predict the ground state to be HS with the value of the gap being proportional to the amount of HF exchange included in the functional.

C. $[\text{Fe}(\text{NCH})_6]^{2+}$

The results for the $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$ and $[\text{Fe}(\text{NH}_3)_6]^{2+}$ ions demonstrate that, although a good choice of basis set might be important to predict accurate molecular structures, the estimated values of the adiabatic energy gaps depend mainly on the functional used. Indeed, for a given functional, two adiabatic gaps obtained with two different basis sets differ at most by a few tens of meV. This has to be compared with the differences in the values predicted by different functionals, which can be of several hundreds meV. For the ion $[\text{Fe}(\text{NCH})_6]^{2+}$, we have then decided to compare only calculations performed using the basis set C, which typically gives us the lowest energy.

$[\text{Fe}(\text{NCH})_6]^{2+}$ has perfect octahedral symmetry in the LS state. In contrast, the structure of the HS state is predicted to have D_{4h} symmetry by B3LYP and PBE0 and C_i symmetry by BP86 and Becke-HH. Table V displays the Fe-ligand bond-lengths and Table VI displays the values of adiabatic energy gap calculated with each functional. Once again these can be ordered according to the series (7). We find that the total energy of the LS state is at least 1 eV lower than that of the HS state for both the LDA and BP86. In contrast PBE0 and B3LYP return the HS state as the most stable, but the absolute value of ΔE^{adia} is only a few hundreds meV (note that

TABLE V. Bond-lengths of $[\text{Fe}(\text{NCH})_6]^{2+}$ in the HS and LS state as calculated with various functionals and for the basis set C.

Functional	Basis set	d_{LS} (Å)	d_{HS} (Å)
LDA	C	1.854	2.066, 2.067, 2.11
BP86	C	1.917	2.171, 2.171, 2.155
B3LYP	C	1.974	2.206, 2.201, 2.201
PBE0	C	1.950	2.194, 2.181, 2.181
HH	C	1.990	2.20, 2.196, 2.196

TABLE VI. Adiabatic energy gap, ΔE^{adia} , for the cation $[\text{Fe}(\text{NCH})_6]^{2+}$ calculated with various functionals and the basis set C.

Functional	Basis set	ΔE^{adia} (cm^{-1})	ΔE^{adia} (eV)
LDA	C	19126.3	2.37135
BP86	C	8410.89	1.04282
B3LYP	C	-1667.48	-0.20674
PBE0	C	-3544.58	-0.43947
HH	C	-12029.62	-1.49148

our B3LYP adiabatic energy gap is consistent with that calculated by Bolvin⁶⁷). Finally, the Becke-HH predicts $\Delta E^{\text{adia}} \approx -1.5$ eV.

D. $[\text{Fe}(\text{CO})_6]^{2+}$

The $[\text{Fe}(\text{CO})_6]^{2+}$ ion has perfect octahedral symmetry in the LS state. This is then reduced to D_{4h} in the HS (the metal-ligand bond-lengths are listed in Table VII). The calculated adiabatic energy gaps are displayed in Table VIII). Again LDA and BP86 are found to (massively) over-stabilize the LS state and the adiabatic energy gap turns out unrealistically large.

At variance to the previous cases, PBE0 and B3LYP return now an almost identical adiabatic energy gaps. In fact, the B3LYP calculated ΔE^{adia} is about 30 meV smaller than the PBE0 one and, therefore, the trend observed through the series in Eq. (7) is not respected. As we will discuss in detail in Sec. V, this result might be related to the fact that the energetic of $[\text{Fe}(\text{CO})_6]^{2+}$ depends largely on the correlation part of the functionals as well as on the exchange part. Finally, we observe that the Becke-HH functional incorrectly predict a HS ground state, meaning that this includes a too large fraction of HF exchange to account accurately for the electronic structure of this ion.

E. Zero-point phononic energies

So far we have focused only on the adiabatic energy gaps. However, the expression for the internal energy difference, Eq. (2), contains also a contribution coming from the phonon zero point energies. Table IX displays ΔE^{ZPE} , calculated by using the various functionals (the results are shown only for the basis set C). ΔE^{ZPE} is found to be always negative (i.e., the zero point energy of the HS state is lower than that of the LS one) reflecting the weaker Fe-ligand bond of the HS

TABLE VII. Bond-lengths of $[\text{Fe}(\text{CO})_6]^{2+}$ in the HS and LS state, as calculated with various functionals and for the basis set C.

Functional	Basis set	d_{LS} (Å)	d_{HS} (Å)
LDA	C	1.848	2.199, 2.172, 2.123
BP86	C	1.900	2.226, 2.331
B3LYP	C	1.948	2.307, 2.367
PBE0	C	1.915	2.276, 2.345
HH	C	1.915	2.322, 2.329, 2.366

TABLE VIII. Adiabatic energy gap, ΔE^{adia} , for the cation $[\text{Fe}(\text{CO})_6]^{2+}$ calculated with various functionals and the basis set C.

Functional	Basis set	ΔE^{adia} (cm^{-1})	ΔE^{adia} (eV)
LDA	C	41 148	5.1017
BP86	C	27 575	3.4189
B3LYP	C	10 656	1.32126
PBE0	C	10 888	1.3501
HH	C	5232	-0.6488

configuration. Corrections to the total energy of the two states then always tend to stabilize the HS.

In contrast to the adiabatic energy gap, ΔE^{ZPE} is found to be almost functional independent. Indeed, for a given system, the difference between two values of ΔE^{ZPE} obtained with two different functionals, are never larger than 15 meV. This demonstrates that the curvature of the PESs is usually very well reproduced by every functional. Therefore, the spread in the predicted values of ΔE^{adia} must arise from the relative shift of the PES of one spin state with respect to that of the other. This observation is consistent with the results by Zein *et al.*,¹⁹ which indicates that the DOGs, defined by Eq. (5), do not depend on the choice of functional.

V. DMC RESULTS AND DISCUSSION

Tables X–XIII display the DMC total energies⁶⁸ of the four ions $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$, $[\text{Fe}(\text{NH}_3)_6]^{2+}$, $[\text{Fe}(\text{NCH})_6]^{2+}$, and $[\text{Fe}(\text{CO})_6]^{2+}$ in both the HS and LS states. The molecular geometries were obtained by DFT optimization (both the functional and the basis set used are indicated in the first column on the left-hand side). Unfortunately, in most cases the DMC energies have a statistical error not small enough to firmly establish which combination of functional/basis set returns the lowest energy structure of a given complex (only the LDA molecular structures have systematically higher energies, but

TABLE IX. Energy difference between the phonon zero point energy of the HS and LS state calculated with the various functionals employed in this work (only results for the basis set C are shown).

System	Functional	ΔE^{ZPE} (eV)	ΔE^{ZPE} (cm^{-1})
$[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$	BP86	-0.08195	-661
$[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$	B3LYP	-0.09079	-732
$[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$	PBE0	-0.09308	-750
$[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$	HH	-0.10272	-828
$[\text{Fe}(\text{NH}_3)_6]^{2+}$	BP86	-0.17413	-1404
$[\text{Fe}(\text{NH}_3)_6]^{2+}$	B3LYP	-0.16099	-1298
$[\text{Fe}(\text{NH}_3)_6]^{2+}$	PBE0	-0.17636	-1422
$[\text{Fe}(\text{NH}_3)_6]^{2+}$	HH	-0.16071	-0.16071
$[\text{Fe}(\text{NCH})_6]^{2+}$	BP86	-0.16593	-1338
$[\text{Fe}(\text{NCH})_6]^{2+}$	B3LYP	-0.15367	-1239
$[\text{Fe}(\text{NCH})_6]^{2+}$	PBE0	-0.13846	-1116
$[\text{Fe}(\text{NCH})_6]^{2+}$	HH	-0.14487	-1168
$[\text{Fe}(\text{CO})_6]^{2+}$	BP86	-0.20800	-1677
$[\text{Fe}(\text{CO})_6]^{2+}$	B3LYP	-0.19894	-1604
$[\text{Fe}(\text{CO})_6]^{2+}$	PBE0	-0.21157	-1706
$[\text{Fe}(\text{CO})_6]^{2+}$	HH	0.21799	-1758

TABLE X. DMC total energy for the LS state, the HS state, and the adiabatic energy gap of the $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$ ion. The molecular structures were optimized by DFT using the various functionals and basis sets listed in the first column. The time-steps chosen for the DMC simulation are also indicated. Differences in energy are well converged for $\Delta\tau = 0.005$ a.u.

Details geom. opt.	$\Delta\tau$ (a.u.)	E_{LS} (eV)	E_{HS} (eV)	ΔE^{adia} (eV)
BP86 (basis C)	0.0125	-6127.211(9)	-6129.720(8)	-2.51(1)
BP86 (basis C)	0.005	-6127.218(9)	-6129.90(2)	-2.65(1)
BP86 (basis C)	0.001	-6127.54(9)	-6130.19(4)	-2.65(9)
B3LYP (basis C)	0.0125	-6127.09(2)	-6129.74(1)	-2.65(2)
B3LYP (basis C)	0.005	-6127.36(1)	-6129.89(2)	-2.54(1)
B3LYP (basis C)	0.002	-6127.44(3)	-6130.01(2)	-2.57(4)
B3LYP (basis C)	0.001	-6127.5(1)	-6130.10(2)	-2.6(1)
PBE0 (basis C)	0.125	-6127.220(9)	-6129.804(8)	2.58(1)
PBE0 (basis C)	0.005	-6127.44(2)	-6129.94(2)	-2.50(3)
PBE0 (basis C)	0.001	-6127.66(6)	-6130.18(4)	-2.52(7)

this is not surprising since the analysis of the phonon modes revealed that these structures are not even associated to a stable minimum of the LDA total energy).

In contrast, the adiabatic energy gaps are calculated with great confidence and they are listed in the right-most column of Tables X–XIII. An analysis of these results can be carried out by looking at Fig. 4, where we present ΔE^{adia} calculated with both DFT and DMC for all the four ions. The systematic up-shift of the LDA and BP86 values with respect to the DMC ones reflects the massive artificial over-stabilization of the LS state (this shift can be as large as few eV). Notably, LDA and BP86 incorrectly return a LS ground state for the ions $[\text{Fe}(\text{NH}_3)_6]^{2+}$ and $[\text{Fe}(\text{NCH})_6]^{2+}$.

B3LYP and PBE0 provide slightly improved results. Their values for ΔE^{adia} lie systematically below those computed with BP86 and the ground state spin is correctly predicted for all ions. However, unfortunately, the quantitative agreement with DMC is still far from being reached as the

TABLE XI. DMC total energy for the LS state, the HS state, and the adiabatic energy gap of the $[\text{Fe}(\text{NH}_3)_6]^{2+}$ ion. The molecular structures were optimized by DFT using the various functionals and basis sets listed in the first column. The time-steps chosen for the DMC simulation are also indicated. Differences in energy are well converged for $\Delta\tau = 0.005$ a.u.

Details geom. opt.	$\Delta\tau$ (a.u.)	E_{LS} (eV)	E_{HS} (eV)	ΔE^{adia} (eV)
LDA (basis C)	0.0125	-5234.92(1)	-5236.93(1)	-2.01(1)
LDA (basis C)	0.005	-5235.33(2)	-5237.17(1)	-1.84(2)
LDA (basis C)	0.001	-5235.69(5)	-5237.36(5)	-1.67(7)
BP86 (basis C)	0.0125	-5235.56(1)	-5237.162(9)	-1.60(1)
BP86 (basis C)	0.005	-5235.78(1)	-5237.37(1)	-1.58(1)
BP86 (basis C)	0.001	-5235.98(3)	-5237.55(5)	-1.57(5)
B3LYP (basis C)	0.0125	-5235.516(9)	-5237.15(1)	-1.63(1)
B3LYP (basis C)	0.005	-5235.77(1)	-5237.36(1)	-1.59(1)
B3LYP (basis C)	0.001	-5236.01(3)	-5237.59(4)	-1.58(5)
PBE0 (basis B)	0.0125	-5235.60(1)	-5237.21(1)	-1.61(1)
PBE0 (basis B)	0.005	-5235.89(2)	-5237.40(2)	-1.51(2)
PBE0 (basis B)	0.001	-5236.14(3)	-5237.67(9)	-1.53(9)
PBE0 (basis C)	0.0125	-5235.57(1)	-5237.133(8)	-1.56(1)
PBE0 (basis C)	0.005	-5235.88(2)	-5237.37(1)	-1.49(1)
PBE0 (basis C)	0.001	-5236.10(3)	-5237.60(2)	-1.50(4)

TABLE XII. DMC total energy for the LS state, the HS state, and the adiabatic energy gap of the $[\text{Fe}(\text{NCH})_6]^{2+}$ ion. The molecular structures were optimized by DFT using the various functionals and basis sets listed in the first column. The time-steps chosen for the DMC simulation are also indicated. Differences in energy are well converged for $\Delta\tau = 0.0125$ a.u.

Details geom. opt.	$\Delta\tau$ (a.u.)	E_{LS} (eV)	E_{HS} (eV)	ΔE^{adia} (eV)
BP86 (basis C)	0.0125	-5957.57(1)	-5959.30(1)	-1.73(2)
BP86 (basis C)	0.005	-5957.57(2)	-5959.32(2)	-1.75(3)
B3LYP (basis C)	0.0125	-5957.94(1)	-5959.32(1)	-1.38(2)
B3LYP (basis C)	0.005	-5957.96(2)	-5959.33(2)	-1.37(3)
PBE0 (basis C)	0.0125	-5957.94(1)	-5959.291(9)	-1.35(1)
PBE0 (basis C)	0.005	-5957.95(3)	-5959.30(1)	-1.35(3)

PBE0 and the DMC results differ by about 0.6 eV (in the best case). Nevertheless, hybrid functionals calculate correctly the relative ligand-field strength of the three HS ions $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$, $[\text{Fe}(\text{NH}_3)_6]^{2+}$, and $[\text{Fe}(\text{NCH})_6]^{2+}$. In fact, although the B3LYP and PBE0 results appear shifted vertically according to the fraction of HF exchange included in the functional, the relative ΔE^{adia} of two complexes is well reproduced. In contrast, the results for $[\text{Fe}(\text{CO})_6]^{2+}$ do not show the same trend. The PBE0 adiabatic energy gap lie slightly above (by about 30 meV) the B3LYP one, despite the fraction of HF exchange being larger in PBE0 than in B3LYP. This indicates that the exchange and correlation energies have different relative importance for the HS and the LS compounds. In order to better understand this important observation, we have calculated the adiabatic energy gap of $[\text{Fe}(\text{NCH})_6]^{2+}$ and $[\text{Fe}(\text{CO})_6]^{2+}$ after changing the fraction of HF exchange in both B3LYP and PBE0 (see Fig. 5).

On the one hand, PBE0 and B3LYP give very similar results for $[\text{Fe}(\text{NCH})_6]^{2+}$, regardless of the different local-exchange and correlation energy. Therefore, the calculated adiabatic energy gaps depend mainly on the fraction of HF exchange (and this dependence is almost linear). This indicates that the correlation contribution to the total energy is well described with (semi-)local functionals and the failures in predicting ΔE^{adia} could be entirely ascribed to the underestimation of the exchange energy. In addition, by fitting the data, we also conclude that about 50% of HF exchange is required to achieve a fair agreement between the DFT and the DMC adiabatic energy gaps. Hence, the Becke-HH functional is found to provide quite satisfactory results (see Fig. 4).

On the other hand, the B3LYP-calculated ΔE^{adia} of $[\text{Fe}(\text{CO})_6]^{2+}$ is systematically down-shifted with respect to the PBE0 value calculated with the same amount of HF exchange. Therefore, for this ion the results depend drastically

TABLE XIII. DMC total energy for the LS state, the HS state, and the adiabatic energy gap of the $[\text{Fe}(\text{CO})_6]^{2+}$ ion. The molecular structures were optimized by DFT using the functionals and the basis sets listed in the first column. The time-steps chosen for the DMC simulation are also indicated. Differences in energy are well converged for $\Delta\tau = 0.0125$ a.u.

Details DFT geom. opt.	$\Delta\tau$ (a.u.)	E_{LS} (eV)	E_{HS} (eV)	ΔE^{adia} (eV)
B3LYP (basis C)	0.0125	-6850.97(2)	-6850.64(2)	0.33(3)
B3LYP (basis C)	0.005	-6850.82(2)	-6850.45(2)	0.37(3)

TABLE XIV. Adiabatic energy gap, ΔE^{adia} , for the four ions calculated with the OLYP and HCTH407 functionals (the basis set C was used).

System	Functional	Basis set	ΔE^{adia} (cm^{-1})	ΔE^{adia} (eV)
$[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$	OLYP	C	-15953	-1.9780
$[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$	HCTH407	C	-19315	-2.3947
$[\text{Fe}(\text{NH}_3)_6]^{2+}$	OLYP	C	-7338	-0.9099
$[\text{Fe}(\text{NH}_3)_6]^{2+}$	HCTH407	C	-9942	-1.2327
$[\text{Fe}(\text{NCH})_6]^{2+}$	OLYP	C	525	0.06510
$[\text{Fe}(\text{NCH})_6]^{2+}$	HCTH407	C	-3650	-0.4526
$[\text{Fe}(\text{CO})_6]^{2+}$	OLYP	C	21313	2.6425
$[\text{Fe}(\text{CO})_6]^{2+}$	HCTH407	C	17097	2.1198

on the correlation as well as the exchange part of the density functional. In addition, we note that about 30% and 40% of HF exchange, respectively, in B3LYP and PBE0, are required to reproduce the DMC gaps and that the HH functional incorrectly describes $[\text{Fe}(\text{CO})_6]^{2+}$ as a HS ion (see Fig. 4).

The DFT performances for the three HS ions and $[\text{Fe}(\text{CO})_6]^{2+}$ are related to the different nature of their ground state wave-function. In fact, this was found to have a much more pronounced multi-configurational character in $[\text{Fe}(\text{CO})_6]^{2+}$ than in all the other complexes,⁶⁹ reflecting the increase in the covalency of the metal-ligand bonds.⁴⁸⁻⁶⁹

Based on this observation, one can reasonably argue that, for the HS complexes the local-part of the exchange-correlation functional is able to capture most of the correlation energy, while the fraction of HF exchange effectively cures the LDA underestimation of the exchange. Thus, hybrid functionals with “enough” HF exchange are found to systematically return quite satisfactory results. Furthermore, as the failures of standard GGA functionals seem mostly related to the shortcomings in the description of the exchange energy, recent GGA functionals, constructed in order to tackle this issue, can out-perform. For example, the OLYP functional, whose exchange part (OPTX) is parametrized to reproduce the Hartree-Fock exchange for atoms,²⁷ predicts values for the adiabatic energy gap of $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$ and $[\text{Fe}(\text{NH}_3)_6]^{2+}$, which compare well with those computed with either B3LYP or

TABLE XV. The adiabatic energy gap for $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$ calculated by using various wave-function methods (reference to the literature is given in the second column). The values labelled as corr-CASSCF and corr-CASPT denote, respectively, the CASSCF and CASPT values after having applied an empirical correction of the order of 4000 cm^{-1} (see main text). Pierloot and Vancoillie⁶⁷ provide an additional long list of results obtained by using different basis sets, geometries, and symmetries. Here, we report only the value that these authors indicate as the “best.”

Method	Reference	ΔE^{adia} (cm^{-1})	ΔE^{adia} (eV)
CASSCF(6,5)	28	-23 125	-2.86714
CASSCF(12,10)	28	-21 180	-2.62599
corr-CASSCF(12,10)	28	-17 892	-2.21833
CASPT2(6,5)	28	-21 610	-2.6793
CASPT2(12,10)	28	-16 185	-2.00668
corr-CASPT2(12,10)	28	-12 347	-1.53083
SORCI	28	-13 360	-1.65643
CASPT2(10,12)	48	-16 307	-2.02181

TABLE XVI. Adiabatic energy gaps for $[\text{Fe}(\text{NH}_3)_6]^{2+}$ calculated by using various wave-function methods (reference to the literature is given in the second column). The values labelled as corr-CASSCF and corr-CASPT denote, respectively, the CASSCF and CASPT values after having applied the empirical correction of the order of 4000 cm^{-1} (see main text).

Method	Reference	$\Delta E^{\text{adia}} (\text{cm}^{-1})$	$\Delta E^{\text{adia}} (\text{eV})$
CASSCF(12,10)	29	-20 630	-2.55779
corr-CASSCF(12,10)	29	-16 792	-2.08194
CASPT2(12,10)	29	-12 963	-1.60721
corr-CASPT2(12,10)	29	-9125	-1.13136
SORCI	29	-10 390/ - 11 250	-1.2882/ - 1.39482
CASPT2(12,10)	48	-7094	-0.879544

PBE0.⁴⁹ Unfortunately, however, OLYP is not as accurate as the hybrids for predicting geometry optimizations and bond-lengths.⁴⁹ Another such GGA functionals, which was found to perform at the level of the hybrids (if not even better),²⁹ is the HCTH407.⁷⁰ Our own results are then listed in Table XIV. The massive improvement, that OLYP and HCTH407 have achieved, over BP86, is evident in the case of $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$, $[\text{Fe}(\text{NH}_3)_6]^{2+}$, and $[\text{Fe}(\text{NCH})_6]^{2+}$.

In contrast, one can question whether $[\text{Fe}(\text{CO})_6]^{2+}$ can be described at all by the single-determinant picture provided by DFT. In principle, the multiconfigurational nature of a wave-function can be described by GGA functionals. In fact the GGA exchange roughly mimics the non-dynamical correlation (in addition to the proper exchange).⁷¹ In practice, however, no DFT flavour investigated here has proven fairly accurate for the energetic of the ion $[\text{Fe}(\text{CO})_6]^{2+}$.

Finally, we compare our DMC results to those obtained with wave-function based methods. Fouqueau *et al.*^{28,29} carried out several calculations for the adiabatic energy gap of the ions $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$ and $[\text{Fe}(\text{NH}_3)_6]^{2+}$ by using the complete

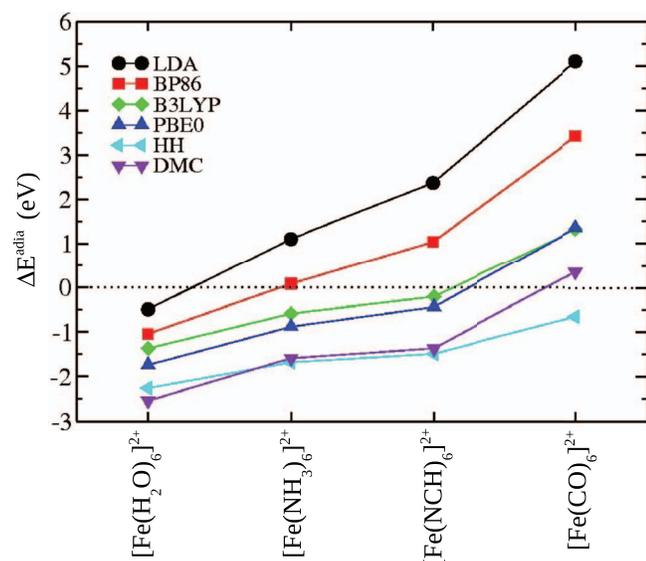


FIG. 4. Adiabatic energy gaps calculated with DFT and DMC. The DFT results were obtained with the functionals indicated in the legend and the basis set C. The DMC results were obtained for the structures optimized with B3LYP (basis set C) and with time-steps $\Delta\tau = 0.005$ a.u. (the error bars are smaller than the symbols).

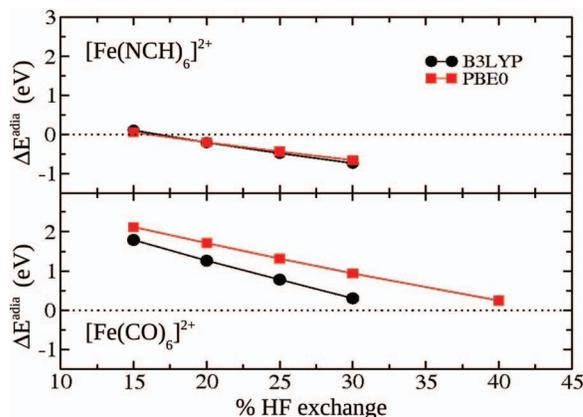


FIG. 5. Adiabatic energy gaps versus the fraction of HF exchange included in the hybrid functionals B3LYP and PBE0 for $[\text{Fe}(\text{NCH})_6]^{2+}$ (upper panel) and $[\text{Fe}(\text{CO})_6]^{2+}$ (lower panel). The basis set C was used.

active space self-consistent field (CASSCF) method with second order perturbative corrections (CASPT2). Some of the results are summarized in Tables XV and XVI. As observed by the authors themselves and in Ref. 48, these calculations suffer the drawback of having been carried out with an insufficient Fe basis set. As such they are affected by a systematic error, which can be estimated by considering the $^5D - ^1I$ splitting of the free Fe^{2+} ion. An empirical correction of the order of 4000 cm^{-1} was then introduced. In the same works, results obtained by spectroscopy-oriented configuration-interaction (SORCI) were also reported. These were stated not to require any empirical correction. A second set of results is provided by Pierloot and Vancoillie,⁴⁸ who performed calculations with basis sets of larger size. For $[\text{Fe}(\text{NH}_3)_6]^{2+}$, their best CASPT2 adiabatic gap agrees fairly well with the corrected-CASPT2 and SORCI results. However, for $[\text{Fe}(\text{H}_2\text{O})_6]^{2+}$, they found a significantly larger (in absolute value) ΔE^{adia} .

By analyzing the data in Tables XV and XVI, we note that the adiabatic energy gaps calculated with CASPT2 by Fouqueau *et al.*^{28,29} agree fairly well with our DMC ones (in particular for $[\text{Fe}(\text{NH}_3)_6]^{2+}$). In contrast, the empirical corrections worsen the agreement and the SORCI results do not agree quantitatively with ours. Although we have not achieved yet a complete understanding of these differences, we argue that they may originate from the large dependence of the CASSCF/CASPT2 results on the basis sets and on the orbitals included in the active space. Furthermore, wave-function based methods do not describe dynamic electronic correlations (although partial corrections are provided by the second order perturbation theory). The DMC energies, in turns, might depend on the choice of the trial wave-function introduced to impose the fixed-node approximation.¹⁵ As described before, the use of the Jastrow factor in the trial-wave functions makes DMC able to retrieve a large fraction of the dynamic correlation energy. In contrast, a very accurate estimate of the static correlation energy would require multi-determinant trial wave-functions, not considered in the present work. This may be quite important for the $[\text{Fe}(\text{CO})_6]^{2+}$ ion, while it may be less relevant for the other three molecules. Furthermore, it was found that, in transition metal oxide molecules and solids,⁵⁰ the DMC energy depends on the degree of *p-d* hybridization

displayed by the single-particle orbitals, which enter in the Slater determinant of the trial wave-function. This might hold even for the molecules investigated in this work. A thorough analysis on these interesting and complex issues is currently on the way and it will be discussed in future works.

VI. CONCLUSIONS

In this work, we have assessed the performances of several popular exchange-correlation functionals in describing various Fe^{2+} complexes. As DFT results cannot be easily related to experiments (at least without accounting for environmental and finite-temperature effects), we have performed accurate DMC calculations, which provide a solid theoretical benchmark for the theory. The DFT and DMC results, both obtained within the theoretical framework of the adiabatic approximation, could be then directly compared.

The LDA and the standard GGA functionals drastically over-stabilize the LS state. Although the accuracy of the DFT calculations increases when hybrid functionals are employed, the most popular ones, B3LYP and PBE0, provide results, which are still quantitatively unsatisfactory when compared against the DMC estimates. In the case of HS ions, a fair agreement between the DFT and the DMC adiabatic gaps is achieved only by using about 50% of HF exchange. In contrast, a lower fraction of HF exchange (between 30% and 40%) is required for $[\text{Fe}(\text{CO})_6]^{2+}$. This difference might be related to the diverse nature of the ground state wave-function for the HS and LS ions. Therefore, unfortunately, we have to conclude that there is not yet a “universal” functional able to correctly describe the energetics of every Fe^{2+} complex. In addition, we remark that our study addressed only the calculation of the adiabatic energy gap. One then can still have the situation where the inclusion of a large fraction of HF exchange in a density functional improves the description of the relative energy of the HS and LS state, but it deteriorates the accuracy of B3LYP and PBE0 in describing other molecular properties. This is a very important issue, which deserves further studies.

Finally, by analyzing zero-point phonon energies, we have demonstrated that the shape of the PESs is well described by every functional considered. Therefore, as already pointed out by Zein *et al.*,¹⁹ the failures of DFT in calculating the adiabatic energy gaps must be ascribed to a shift of the PES of the LS state with respect to that of the HS state.

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